# COMPLEX FORMATION AND KINEMATIC COUPLING: $\mathbf{C}_{2} \mathbf{H}_{\mathbf{4}} \cdot \mathrm{PtCl}_{3}{ }^{*}$ 

Roman Řeřicha ${ }^{a}$, Björg N. Cyvin ${ }^{b}$, Jon Brunvoll ${ }^{b}$ and Sven J. Cyvin ${ }^{b}$<br>${ }^{a}$ Institute of Chemical Process Fundamentals, Czechoslovak Academy of Sciences, 16502 Prague-Suchdol, Czechoslovakia and<br>${ }^{b}$ Division of Physical Chemistry,<br>University of Trondheim, N-7034 Trondheim - NTH, Norway

Received January 26th, 1988
Accepted February 8th, 1988

Dedicated to Dr R. Zahradnik on the occasion of his 60th birthday.

Normal coordinate analyses including calculations of $P E D$ 's were performed for $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ system modelling Zeise's anion, $\left[\left(\mathrm{C}_{2} \mathrm{H}_{4}\right) \mathrm{PtCl}_{3}\right]^{-}$. The wedgewise distorsion of the $\mathrm{C}_{2} \mathrm{H}_{4}$ ligand known from the crystallographic data for Zeise's salt, was taken into account. Under these circumstances it was found that the kinematic couplings between the internal ligand and complex framework vibrational modes are rather small. The reliability of some existing assignments of the fundamental frequencies of Zeise's anion is discussed.

Zeise's salt, $\mathrm{K}\left[\left(\mathrm{C}_{2} \mathrm{H}_{4}\right) \mathrm{PtCl}_{3}\right] \cdot \mathrm{H}_{2} \mathrm{O}$, was the first organometallic compound synthetized by $\operatorname{man}^{1}$ around 1830 . Since 1950's, a huge amount of experimental data on this complex has been accumulated by various physical methods, however, sometimes with confficting interpretations of the data. Studies of this stable complex can give information about a catalytically active particle $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PdCl}_{3}$ predicted by analogy, the structure of which is not known for understandable reasons (cf. ref. ${ }^{2}$ ).

The electronic problem of the alkene-Pt(II) coordination bond can be attacked indirectly, using vibrational spectroscopy. On comparing the characteristic vibrational frequencies of the free and coordinated molecule of ethene, frequency shifts can be considered as an indirect measure of electronic changes of the ligand. However, an unambiguous assignment of the observed spectral bands of the complex to proper characteristic vibrations of $\mathrm{C}_{2} \mathrm{H}_{4}$ after its coordination remains to be a problem.

Moreover, it has been shown by calculations, e.g. ref. ${ }^{3}$, that some frequency shifts can result from other reasons than from changes in the potential energy of the $\mathrm{C}_{2} \mathrm{H}_{4}$ ligand. This phenomenon, entitled as the kinematic coupling effect ${ }^{3}$, results from non-vanishing off-diagonal inverse kinetic energy matrix elements which couple

* Part LXXII in the series Catalysis by Metal Complexes; Part LXXI: React. Kinet. Catal. Lett. 31, 309 (1986).
the internal ligand and complex framework modes (CFM). This has been demonstrated ${ }^{3}$ on a hypothetical molecule $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{Tl} \cdot \mathrm{H}_{2} \mathrm{O}$, which is interesting from the point of view of catalysis by trivalent thallium, where the molecule of ethene has been considered in the planar configuration of nuclei.
On the other hand, a wedgewise distorsion of coordinated $\mathrm{C}_{2} \mathrm{H}_{4}$ in stable Zeise's anion is known experimentally since 1970 's from neutron ${ }^{4}$ and X-ray ${ }^{5}$ diffraction data. In our model calculations for $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ complexed particle, eighteen diagonal F matrix elements (twelve for $\mathrm{C}_{2} \mathrm{H}_{4}$ and six for $\mathrm{PtCl}_{3}$ ) were kept constant and the six remaining ones (the CFM's) were varied in order to observe trends in eighteen corresponding frequencies and in potential energy distributions (PED's).

The aim of this work was to compare the kinematic coupling effects for planar $\mathrm{C}_{2} \mathrm{H}_{4}$ in the thallic complex on one hand and the wedgewise distorted $\mathrm{C}_{2} \mathrm{H}_{4}$ complexed with the $\mathrm{Pt}(\mathrm{II})$ species on the other.

## THEORETICAL

Calculations. All the calculations were performed on a UNIVAC computer of the Norwegian Institute of Technology in Trondheim, using the standard programs for normal coordinate analyses (cf. ref. ${ }^{3}$ ).

Symmetry coordinates of molecular vibrations. The idealized model of Zeise's anion belongs to the symmetry $C_{2 v}$. When referring to the coordinate axes defined in Fig. 1, the normal vibrations are distributed into the various symmetry species as:

$$
\Gamma_{\text {vib }}=8 a_{1}+4 a_{2}+6 b_{1}+6 b_{2} .
$$



Fig. 1
The idealized geometry of Zeise's anion with the Kohlrausch-Herzberg orientation of $\boldsymbol{x}^{\prime}$, $y^{\prime}$, and $z \equiv z^{\prime}$ axes and $C_{2 v}$ symmetry. The bond lengths used (in pm): CH (108.7), CC (137.4), Pt-midpoint CC (203), $\mathrm{Cl}-\mathrm{Pt}$ ( 234 for the atoms $9-10$ and 230 for both $7-10$ and $8-10) ; \mathrm{HCH}\left(117^{\circ}\right)$, wedge: angle of $\mathrm{CH}_{2}$ planes ( $145^{\circ}$ )

[^0]These modes of vibration may be divided into three types: the ethene vibrations, those of the $\mathrm{PtCl}_{3}$ fragment, and the couplings between these two parts of the complex ( $C F M$ 's).

For $\mathrm{C}_{2} \mathrm{H}_{4}$ we have considered both the planar configuration of nuclei (symmetry $D_{2 h}$ ), and the wedgewise distorted one which pertains to Zeise's anion in crystals (symmetry $C_{2 v}$; cf. Fig. 1). The correlation of the normal vibrations of these two configurations is given below, where the species of $D_{2 h}$ are denoted in parentheses (only here) by capital letters.

$$
\begin{aligned}
& \Gamma_{\mathrm{vib}}\left(\mathrm{C}_{2} \mathrm{H}_{4}\right)=4 a_{1}\left(3 A_{g}+B_{1 u}\right)+3 a_{2}\left(2 B_{1 g}+A_{u}\right)+ \\
& \quad+3 b_{1}\left(B_{2 g}+2 B_{3 u}\right)+2 b_{2}\left(2 B_{2 u}\right)
\end{aligned}
$$

In a preliminary set of calculations we used the symmetry coordinates for $\mathrm{C}_{2} \mathrm{H}_{4}$ from Cyvin and Cyvin ${ }^{6,7}$. In the next round we wished to utilize the presumably best force field for free planar ethene due to Duncan and Hamilton ${ }^{8}$, where 116 experimental frequencies were used in the refinement of the force constants. Hence we executed the slight changes of the in-plane symmetry coordinates in order to conform the definitions of Duncan et al. ${ }^{8}$, adopted from Kuchitsu et al. ${ }^{9}$. These definitions imply unscaled bending coordinates.

For the part associated with the $\mathrm{PtCl}_{3}$ fragment (cf. ref. ${ }^{10}$ ), which has the same symmetry $\left(C_{2 v}\right)$ as the whole complex, one obtains the following symmetry species :

$$
\Gamma_{\mathrm{vib}}\left(\mathrm{PtCl}_{3}\right)=3 a_{1}+b_{1}+2 b_{2} .
$$

For the whole complex there remain six vibrational modes which couple the internal ligand and complex framework modes (CFM). They are distributed into the symmetry species according to:

$$
\Gamma_{\mathrm{vib}}(C F M)=a_{1}+a_{2}+2 b_{1}+2 b_{2} .
$$

For the sake of simplicity of their description, these CFM's can be characterized as the original translational $\left(T_{x}^{\prime}, T_{y}^{\prime}, T_{z}\right)$ and rotational $\left(R_{x}^{\prime}, R_{y}^{\prime}, R_{z}\right)$ degrees of freedom of the $\mathrm{C}_{2} \mathrm{H}_{4}$ molecule with respect to the original axis system $x^{\prime}, y^{\prime}, z \equiv z^{\prime}$ (hence the primed quantities; cf. Fig. 1). These originally non-genuine vibrational degrees of freedom of the free $\mathrm{C}_{2} \mathrm{H}_{4}$ molecule are hindered in the field of $\mathrm{PtCl}_{3}$, and they become the genuine vibrational degrees of freedom of the whole complex.

For instance, the $T_{z}$ mode represents the Pt-ethene stretching vibration, the other ones are deformation modes with the exception of $R_{y}^{\prime}$, which can be regarded as another stretching mode with either a large force constant (the so-called antisymmetric $\mathrm{PtC}_{2}$ ring stretching) or with a relatively small one (the so-called $\mathrm{C}_{2} \mathrm{H}_{4}$ pivoting or $\mathrm{C}_{2} \mathrm{H}_{4}$ tilting). We have studied both possibilities, the survey of which is given in Table I.

Precise definitions of the six coupling coordinates are given below:

$$
\begin{aligned}
a_{1}: T_{z} & =2^{-1 / 2}\left(u_{1}+u_{2}\right) \\
a_{2}: R_{z} & =\left(\eta_{1}-\eta_{2}+\eta_{3}-\eta_{4}\right) / 2 \\
b_{1}: T_{x}^{\prime} & =2^{-1 / 2}\left(\Theta_{1}-\Theta_{2}\right) \\
R_{y}^{\prime} & =2^{-1 / 2}\left(u_{1}-u_{2}\right) \\
b_{2}: T_{y}^{\prime} & =\left(\eta_{1}-\eta_{2}-\eta_{3}+\eta_{4}\right) / 2 \\
R_{x}^{\prime} & =\left(w_{1}-w_{2}-w_{3}+w_{4}\right) / 2 .
\end{aligned}
$$

Here $u$ designates the PtC stretchings, while the other symbols refer to bendings scaled by factors of the dimension of length in the usual way ${ }^{6,7}$. Specifically, $\eta$ and $\Theta$ pertain to the two types of CPtCl bending, where in the latter case $(\Theta) \mathrm{Cl}$ stands for the unique (axial) Cl atom, and $w$ pertains to HCPt angles.

The wedgewise distorsion of $\mathrm{C}_{2} \mathrm{H}_{4}$ and its complexation with $\mathrm{PtCl}_{3}$ requires a change in the notation of normal vibrations of the planar molecule which, however, will be used in the description of $P E D$. For that reason, a correlation of both notations for planar and wedgewise distorted $\mathrm{C}_{2} \mathrm{H}_{4}$ is presented in Table II.

## RESULTS AND DISCUSSION

The calculated frequencies are given in Table III together with the corresponding PED's. In column 2, the notation for planar $\mathrm{C}_{2} \mathrm{H}_{4}$ is indicated (cf. Table II), in column 3 the frequencies for the free $\mathrm{PtCl}_{3}$ fragment and the free wedgewise distorted $\mathrm{C}_{2} \mathrm{H}_{4}$ are collected.

The data for the $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ complex are presented in column 4 for the case of the

Table I
The six complex framework modes (CFM) and their arbitrarily chosen force constants (in $\mathrm{mdyn} / \AA)^{a}$ for three types of interactions between $\mathrm{C}_{2} \mathrm{H}_{4}$ and $\mathrm{PtCl}_{3}$

| $C F M$ | Interaction |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
|  | weak | moderate | strong |  | stretching $^{b}$ | deformation $^{c}$ |
| Stretching | 0.10 | 0.20 | 1.00 | $R_{y}^{\prime}, T_{z}$ | $T_{z}$ |  |
| Deformation | 0.01 | 0.05 | 0.10 | $R_{z}, T_{y}^{\prime}, R_{x}^{\prime}, T_{x}^{\prime}$ | $R_{z}, T_{y}^{\prime}, R_{x}^{\prime}, T_{x}^{\prime}, R_{y}^{\prime}$ |  |

[^1]strong interaction only (cf. Table I), which sufficiently illustrates the results of this study together with the additional Tables IV and V.

It is clear from these data that kinematic coupling effects in $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ are relatively small, and on comparing them with the case of $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{Tl} \cdot \mathrm{H}_{2} \mathrm{O}$ they are smaller approximately by one order of magnitude. It is best illustrated by the $\mathrm{CH}_{2}$ rocking vibration $\left(b_{2 u}\right)$, the shift of which for $\mathrm{Tl}^{3+}$ was $+157 \mathrm{~cm}^{-1}$ whilst for $\mathrm{Pt}(\mathrm{II})$ is $+24 \mathrm{~cm}^{-1}$ only.

It follows from this fact that occasional large frequency shifts in the vibrational spectra of Zeise's anion must be attributed to changes in the potential energy of $\mathrm{C}_{2} \mathrm{H}_{4}$ after its coordination to $\mathrm{Pt}(\mathrm{II})$. However, as it was stated above the problem is the unambiguous assignment of the observed spectral bands of the complex to proper characteristic vibrations of $\mathrm{C}_{2} \mathrm{H}_{4}$ after its coordination.

This problem can be well illustrated by two examples of different types of coordination of the ethene molecule to $\mathrm{Pt}(0)$ or $\mathrm{Pt}(\mathrm{II})$.

In case of a reaction between $\mathrm{C}_{2} \mathrm{H}_{4}$ and tetrakis(triphenylphosphine) platinum(0) with the coordination number 4 , two phosphine ligands are replaced by one molecule of ethene.

$$
\begin{equation*}
\mathrm{C}_{2} \mathrm{H}_{4}+\left[\mathrm{Pt}\left(\left(\mathrm{C}_{6} \mathrm{H}_{5}\right)_{3} \mathrm{P}\right)_{4}\right] \rightarrow 2\left(\mathrm{C}_{6} \mathrm{H}_{5}\right)_{3} \mathrm{P}+\left[\left(\mathrm{C}_{2} \mathrm{H}_{4}\right) \mathrm{Pt}\left(\left(\mathrm{C}_{6} \mathrm{H}_{5}\right)_{3} \mathrm{P}\right)_{2}\right] \tag{A}
\end{equation*}
$$

## Table II

Correlation between the notations of some normal vibrations for planar and complexed wedgewise distorted $\mathrm{C}_{2} \mathrm{H}_{4}$

| Free planar $\mathrm{C}_{2} \mathrm{H}_{4}{ }^{a}$ | Complexed wedgewise $\mathrm{C}_{2} \mathrm{H}_{4}{ }^{b}$ |
| :--- | :--- |
|  |  |
| $T_{z}\left(b_{1 u}\right)$ | $T_{z}\left(a_{1}\right):$ stretch. $\mathrm{Pt}-\mathrm{C}_{2} \mathrm{H}_{4}$ |
| rock. $\mathrm{CH}_{2}\left(b_{1 g}\right)$ | twist. $\mathrm{CH}_{2}\left(a_{2}\right)$ |
| CC torsion $\left(a_{u}\right):$ twist. $\mathrm{CH}_{2}$ | CC torsion $\left(a_{2}\right):$ rock. $\mathrm{CH}_{2}$ |
| $R_{z}\left(b_{1 g}\right)$ | $R_{z}\left(a_{2}\right):$ hindered $\mathrm{C}_{2} \mathrm{H}_{4}$ rotation |
| $R_{y}^{\prime}\left(b_{2 g}\right)$ | $\mathrm{C}_{2} \mathrm{H}_{4}$ pivoting or tilting $\left(b_{1}\right)^{d}$ |
| $T_{x}^{\prime}\left(b_{3 u}\right)$ | $T_{\mathrm{CC}}\left(b_{1}\right):$ o-o-p def. $\mathrm{C}_{2} \mathrm{H}_{4}-\left(\mathrm{PtCl}_{3}\right)^{e}$ |
| rock. $\mathrm{CH}_{2}\left(b_{2 u}\right)$ | twist. $\mathrm{CH}_{2}\left(b_{2}\right)$ |
| $T_{y}^{\prime}\left(b_{2 u}\right)$ | $\left(b_{2}\right):$ in-plane def. $\mathrm{C}_{2} \mathrm{H}_{4}-\left(\mathrm{PtCl}_{3}\right)^{f}$ |
| $R_{x}^{\prime}\left(b_{3 g}\right)$ | $R_{\mathrm{CC}}\left(b_{2}\right):$ rock. $\mathrm{CH}_{2}{ }^{c}$ |

[^2]The reaction $(A)$ indicates that the molecule of ethene behaves as a bidentate ligand which is accompanied by a marked lengthening of the CC bond from 133 to 143 pm after the coordination ${ }^{11}$ (cf. 147 pm in the cyclic oxirane ${ }^{12} \mathrm{C}_{2} \mathrm{H}_{4} \mathrm{O}$, and 149 pm in thiirane ${ }^{13}$ ). Thus, the assumption of a three-membered platinirane $\mathrm{PtC}_{2}$ ring formation is substantiated in the Eq. (A) and, in this way, the original coordination number 4 is preserved for the central zerovalent Pt atom in the product.

Table III
Frequencies of the free wedgewise distorted $\mathrm{C}_{2} \mathrm{H}_{4}$, free planar $\mathrm{PtCl}_{3}$, and complexed $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ species (in $\mathrm{cm}^{-1}$ ) and $P E D$ 's

| $C_{2}$ | Notation ${ }^{\text {a }}$ | Free moieties (PED) | Complex, strong int. $(P E D)$ |
| :---: | :---: | :---: | :---: |
| $8 a_{1}$ | CH stretch. | 3013 (99+1CC) | 3014 (99+1CC) |
|  | CC stretch. | $1607(63+36 \mathrm{sc})$ | $1611(63+36 \mathrm{sc})$ |
|  | sciss. $\mathrm{CH}_{2}$ | 1350 (69+28 CC + 2 w) | $1351(69+28 \mathrm{CC}+2 \mathrm{w})$ |
|  | wag. $\mathrm{CH}_{2}\left(b_{1 u}\right)$ | 923 (94+6CC) | $931\left(91+7 \mathrm{CC}+2 T_{z}\right)$ |
|  | ClPtCl str. | $312(58+42 \mathrm{PtCl})$ | $324(77+11 \mathrm{PiCl}+$ |
|  |  |  | $\left.11 T_{z}\right)$ |
|  | PiCl stretch. | $331\left(59+41 \mathrm{PiCl}_{2}\right)$ | $\begin{gathered} 303\left(60+23 \mathrm{PtCl}_{2}+\right. \\ \left.16 T_{z}\right) \end{gathered}$ |
|  | i-p def. ClPtCl | 151 (99) | $149\left(99+1 T_{z}\right)$ |
|  | $T_{z}$ (stretch.) |  | $363\left(69+\operatorname{see}^{\text {b }}\right.$ ) |
| $4 a_{2}$ | CH stretch. | 3084 (100) | 3084 (100) |
|  | rock. $\mathrm{CH}_{2}\left(b_{1 g}\right)$ | $1207(95+5$ (w) | $1212(95+5 \mathrm{tw})$ |
|  | twist. $\mathrm{CH}_{2}$ | $1009(95+5$ r) | $1009(95+5 \mathrm{r})$ |
|  | $R_{z}(\text { def. })$ |  | 131 (99) |
| $6 b_{1}$ | CH stretch. | 3013 (100) | $3013(100)$ |
|  | sciss. $\mathrm{CH}_{2}$ | 1429 (98+2 w) | $1430(98+2 w)$ |
|  | wag. $\mathrm{CH}_{2}$ | 929 (98 + 2 sc ) | (see |
|  | o-o-p def. ClPtCl | 149 (100) | Table IV |
|  | $R_{y}^{\prime}$ (str. or def.) |  | and |
|  | $T_{y}^{\prime}$ (def.) |  | Table V |
| $6 b_{2}$ | CH stretch. | 3104 (100) | 3105 (99) |
|  | rock. $\mathrm{CH}_{2}\left(b_{2 u}\right)$ | 824 (100) | $848\left(89+10 R_{x}^{\prime}+1 \mathrm{CH}\right)$ |
|  | ClPtCl stretch. | $319(98+2 \mathrm{P} \mathrm{Cl})$ | $321\left(96+2 \mathrm{PlCl}+2 T_{y}^{\prime}\right)$ |
|  | $\mathrm{i}-\mathrm{p} \text { def. } \mathrm{PtCl}$ | $198\left(99+1 \mathrm{PtCl}_{2}\right)$ | $211\left(65+34 T_{y}^{\prime}\right)$ |
|  | $T_{y}^{\prime}$ (def.) |  | $171(63+33 \mathrm{PtCl}+$ |
|  |  |  | $\left.2 \mathrm{PtCl}_{2}\right)$ |
|  | $R_{x}^{\prime}($ def.$)$ |  | $567(89+10$ r) |

[^3]Table IV
Extension of Table III for three interaction cases: $R_{y}^{\prime}$ is a stretching mode of $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ complex (in $\mathrm{cm}^{-1} ; P E D$ in $\%$ )

| Notation ( $b_{1}$ modes) | Free | Weak | Moderate | Strong |
| :---: | :---: | :---: | :---: | :---: |
| wag. $\mathrm{CH}_{2}\left(b_{2 g}\right)$ | 929 | $931(98+2 \mathrm{sc})$ | $934\left(97+2 \mathrm{sc}+1 R_{y}^{\prime}\right)$ | $955\left(93+6 T_{x}^{\prime}+1 \mathrm{sc}\right)$ |
| o-o-p def. ClPtCl | 149 | $150\left(98+1 T_{x}^{\prime}\right)$ | $159\left(69+17 T_{x}^{\prime}+14 R_{y}^{\prime}\right)$ | $135\left(70+30 T_{x}^{\prime}\right)$ |
| $R_{y}^{\prime}(\text { str. })^{a}$ |  | $99\left(98+1(\right.$ oop $\left.)+1 T_{x}^{\prime}\right)$ | 137 (80 + 19(oop)) | $307\left(93+5 \mathrm{w}+1 T_{x}^{\prime}\right)$ |
| $T_{x}^{\prime}(\text { def. })^{a}$ |  | $51\left(98+1(\mathrm{oop})+1 R_{y}^{\prime}\right)$ | $108\left(83+13(\mathrm{oop})+4 R_{y}^{\prime}\right)$ | $176\left(68+30(\right.$ oop $\left.)+1 R_{y}^{\prime}\right)$ |
| ${ }^{a}$ For the force constants of $R_{y}^{\prime}$ and $T_{x}^{\prime} \mathrm{cf}$. Table I. |  |  |  |  |
| Table V |  |  |  |  |
| Extension of Table III for three interaction cases: $R_{y}^{\prime}$ is a deformation mode of $\mathrm{C}_{2} \mathrm{H}_{4} \cdot \mathrm{PtCl}_{3}$ complex (in $\mathrm{cm}^{-1} ; P E D$ in $\%$ ) |  |  |  |  |
| Notation ( $b_{1}$ modes) | Free | Weak | Moderate | Strong |
| wag. $\mathrm{CH}_{2}\left(b_{2 g}\right)$ | 929 | 929 (98+2 sc) | $930(98+2 \mathrm{sc})$ | 932 (98+2 sc) |
| $\mathrm{o}-\mathrm{o}-\mathrm{p}$ def. ClPtCl | 149 | $150\left(99+1 T_{x}^{\prime}\right)$ | $156\left(84+16 T_{x}^{\prime}\right)$ | $135\left(71+28 T_{x}^{\prime}\right)$ |
| $R_{y}^{\prime}(\text { def. })^{a}$ |  | $31\left(99+1 T_{x}^{\prime}\right)$ | $69\left(98+1 T_{x}^{\prime}\right)$ | $97\left(98+1 T_{x}^{\prime}\right)$ |
| $T_{x}^{\prime}(\text { def. })^{a}$ |  | $51\left(98+1(\mathrm{oop})+1 R_{y}^{\prime}\right)$ | $110\left(82+16(\mathrm{opp})+1 R_{y}^{\prime}\right)$ | $180\left(70+29(\right.$ oop $\left.)+1 R_{y}^{\prime}\right)$ |

${ }^{a}$ For the force constants of $R_{y}^{\prime}$ and $T_{x}^{\prime} \mathrm{cf}$. Table I.

On the other hand, in the reaction between ethene and the $\mathrm{PtCl}_{4}^{2-}$ ion (in acidified aqueous solutions) only one chloride ligand is replaced by one molecule of ethene to form Zeise's anion, $\left[\left(\mathrm{C}_{2} \mathrm{H}_{4}\right) \mathrm{PtCl}_{3}\right]^{-}$. In this case the molecule of ethene behaves as a monodentate ligand preserving in $\mathrm{Pt}(\mathrm{II})$ species the coordination number 4 which is usual for the stable compounds of this type. It is accompanied by a relatively small lengthening of the CC bond of ethene from 133 to 137 pm after its coordination ${ }^{4.5}$.
On the basis of these facts we believe that the assumption of the $\mathrm{PtC}_{2}$ ring formation for $\operatorname{Pt}(0)$ in Eq. $(A)$ is justified, and thus the $R_{y}^{\prime}$ mode may be the antisymmetric $\mathrm{PtC}_{2}$ ring stretching vibration with a large force constant, generally $R_{y}^{\prime}>T_{z}$.

On the other hand, the same assumption of the $\mathrm{PtC}_{2}$ ring formation in Zeise's anion is not acceptable in view of the small increase in the CC bond length. Moreover, this assumption would mean an unusual coordination number 5 for this stable $\mathrm{Pt}(\mathrm{II})$ species. In other words, the $R_{y}^{\prime}$ mode in this case should be a mode with a relatively low frequency (the so-called $\mathrm{C}_{2} \mathrm{H}_{4}$ pivoting or tilting; generally $T_{z}>R_{y}^{\prime}$; cf. refs ${ }^{14,15}$ ).
This is a basic difference between two kinds of the molecular ethene ligand coordinated to two different stages of the Pt species, $\mathrm{Pt}(0)$ and $\mathrm{Pt}(\mathrm{II})$. Further studies of this problem seem to be warranted.

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[^0]:    Collection Czechoslovak Chem. Commun. (Vol. 53) (1988)

[^1]:    ${ }^{a} 1 \mathrm{mdyn} / \AA=10^{-2} \mathrm{~N} / \mathrm{m} ;{ }^{b}$ the antisymmetric $\mathrm{PtC}_{2}$ ring stretching (i.e., the coordination number 5); ${ }^{c}$ i.e., the $\mathrm{C}_{2} \mathrm{H}_{4}$ pivoting or $\mathrm{C}_{2} \mathrm{H}_{4}$ tilting (and the coordination number 4).

[^2]:    ${ }^{a}$ Axes: $x^{\prime}$ along the $\mathrm{C}=\mathrm{C}$ bond, $y^{\prime}, z=z^{\prime}$ (perpendicular to the original molecular plane, i.e., the symmetry species according to Kohlrausch ${ }^{16}$ and $\operatorname{Herzberg}^{17}$; ${ }^{b}$ axes: $x$ (parallel to $x^{\prime}$ ), $y$ (parallel to $y^{\prime}$ ), $z=z^{\prime}$, cf. Fig. $1 ;{ }^{c}$ also: $\mathrm{PtC}_{2}$ twisting (ref. ${ }^{18}$ ) or $\mathrm{C}_{2} \mathrm{H}_{4}$ torsion (ref. ${ }^{19}$ ); ${ }^{d}$ alternatively: the antisymmetric $\mathrm{PtC}_{2}$ ring stretching (ref. ${ }^{18}$ ); ${ }^{e}$ also: $\mathrm{PtC}_{2}$ rocking (ref. ${ }^{18}$ ); ${ }^{\rho}$ also: $\mathrm{PiC}_{2}$ wagging (ref. ${ }^{18}$ ); ${ }^{g}$ also: a barrelling motion (ref. ${ }^{19}$ ).

[^3]:    ${ }^{a}$ Notation for planar $\mathrm{C}_{2} \mathrm{H}_{4}$; for changes cf. Table II; ${ }^{b} 69 T_{s}+27 v(\mathrm{Pt}-\mathrm{Cl})+2$ wag. $\mathrm{CH}_{2}+$ $1(\mathrm{i}-\mathrm{p})$ def. $\mathrm{PtCl}_{2}+1 \nu\left(\mathrm{PtCl}_{2}\right)$.

